

# PHYS 314

## HOMEWORK #7

### Solutions

1. Consider the problem, done in class and in the text, of finding the gravitational force due to a spherical shell at points exterior to the shell, interior to the shell, and in the shell. In each case, derive the limits for the  $dr$  integral and verify that the text is using the correct values. Verify the integration resulting in eq. 5.21

**Solution :** Let's begin with eq. (5.16) in the text :

$$\Phi = \frac{-2\pi\rho G}{R} \int_b^a r' dr' \int_{r_{\min}}^{r_{\max}} dr \quad (1)$$

where  $r$ ,  $r'$  and  $R$  are related by the law of cosines:

$$r^2 = r'^2 + R^2 - 2r'R \cos \theta$$

a) For the case where  $R > a$ ,  $\theta$  will vary from 0 to  $\pi$ ; when  $\theta = 0$ , the law of cosines gives us

$$r^2 = (r')^2 + R^2 - 2r'R = (R - r')^2 \Rightarrow r = R - r'$$

In this case, we know that  $R > r'$ , so we know that  $R - r'$  must be positive and this will be the lower limit of the  $dr$  integral. When  $\theta = \pi$  we have

$$r^2 = (r')^2 + R^2 - 2r'R(-1) = (r' + R)^2 \Rightarrow r = R + r'$$

This is our upper limit on the  $dr$  integral, so that equation (1) becomes:

$$\begin{aligned} \Phi &= \frac{-2\pi\rho G}{R} \int_b^a r' dr' [R + r' - (R - r')] = \frac{-2\pi\rho G}{R} \int_b^a r' dr' (2r') \\ &= \frac{-4\pi\rho G}{R} \int_b^a (r')^2 dr' \end{aligned}$$

from which eq. (5.17) follows simply.

b) When  $R < b$ , we know that  $r' > R$ , so that when  $\theta = 0$ ,  $r = (r' - R)$  and  $r = (r' + R)$  when  $\theta = \pi$ .

With these limits, we have :

$$\Phi = \frac{-2\pi\rho G}{R} \int_b^a r' dr' \int_{r_{\min}}^{r_{\max}} dr = \frac{-2\pi\rho G}{R} \int_b^a r' dr' \int_{r'-R}^{r'+R} dr$$

the last integral yields  $(r'+R)-(r'-R) = 2R$  which leads to:

$$\Phi = \frac{-2\pi\rho G}{R} \int_b^a r' dr' \int_{r_{\min}}^{r_{\max}} dr = \frac{-2\pi\rho G}{R} \int_b^a r' dr' (2R) = -2\pi G\rho (a^2 - b^2)$$

which significantly does not bear any R dependence.

c) Finally, when R lies in the shell, we obtain two integrals; the first integral sums contributions to  $\Phi$  from points  $b < R < r'$ , and the second integral sums contributions from points  $r' < R < a$ , therefore,  $\Phi$  is the sum of the two integrals :

$$\begin{aligned} \Phi &= \frac{-2\pi\rho G}{R} \left[ \int_R^a r' dr' \int_{r'-R}^{r'+R} dr + \int_b^R r' dr' \int_{R-r'}^{R+r'} dr \right] \\ &= \frac{-2\pi\rho G}{R} \left[ \int_R^a 2Rr' dr' + \int_b^R 2r'r' dr' \right] = -2\pi\rho G (a^2 - R^2) - \frac{4\pi\rho G}{3R} (R^3 - b^3) \end{aligned}$$

and these combine algebraically to eq. 5.21.

2. Verify eqs. 5.22 in the text.

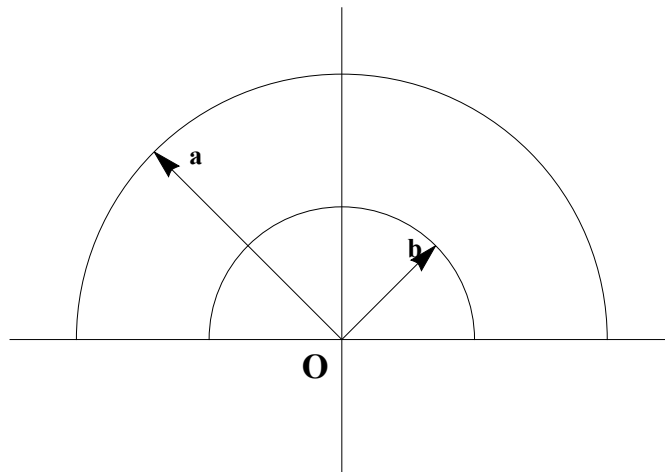
**Solution** : Gravity is a conservative force, and therefore is derived by taking the gradient of its appropriate scalar potential. If we take  $-\nabla\Phi$  for each of the cases above, we find :

$$\text{a) } g(R > a) = -\frac{d}{dR} \left( \frac{-GM}{R} \right) = -\frac{GM}{R^2}$$

$$\text{b) } g(R < b) = \frac{-d}{dR} (-2\pi\rho G (a^2 - b^2)) = 0 \text{ since there is no R dependence}$$

$$\begin{aligned} \text{c) } g(b < R < a) &= \\ \frac{-d}{dR} \left( -4\pi\rho G \left( \frac{a^2}{2} - \frac{b^3}{3R} - \frac{R^2}{6} \right) \right) &= -4\pi\rho G \left( \frac{b^3}{3R^2} - \frac{R}{3} \right) = \frac{-4}{3} \pi\rho G \left( \frac{b^3}{R^2} - \frac{R}{3} \right) \end{aligned}$$

3. A uniform plate has its boundary consisting of two concentric half circles of radii a and b as shown below. Find the force of attraction on a test mass located at the origin (at point O).



**Solution** : We consider the force of attraction of an element of mass within the concentric shells on a test mass  $m$  at the origin. We can write :

$$d\mathbf{F} = -\frac{G m dM}{r^2} \hat{\mathbf{r}}$$

where  $m$  is the mass of the test  $m$ ,  $dM$  is the mass of the element of area, and  $r$  is the distance between the origin and  $dM$ . Since the shell is uniform, we know that  $dM = \rho dA$  and in polar coordinates, we can write  $dA = r dr d\theta$ . (In this problem, we will use  $r$  to denote radial distance to avoid confusion with  $\rho$  for density.) Our expression for  $dF$  becomes :

$$d\mathbf{F} = -\frac{G m \rho r dr d\theta}{r^2} = -\frac{G m \rho dr d\theta}{r} \hat{\mathbf{r}}$$

We can use symmetry arguments to show that the  $x$  components of the force will cancel, leaving us with only a  $y$  component of force. The incremental component of  $dF$  in the  $y$  direction is:

$$dF_y = -\frac{G m \rho dr d\theta \sin \theta}{r}$$

where  $\theta$  is measured counterclockwise from the positive  $x$  axis. Our limits of integration in  $r$  and  $\theta$  give us the following integral :

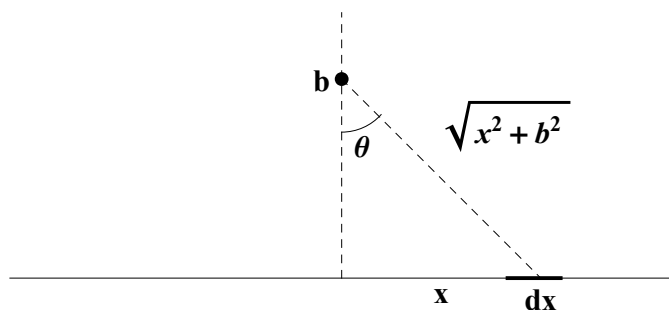
$$F_y = -G m \rho \int_b^a \frac{dr}{r} \int_0^\pi \sin \theta d\theta = -2 G m \rho \int_b^a \frac{dr}{r} = -2 G m \rho \ln \left( \frac{a}{b} \right)$$

you can eliminate  $\rho$  by writing:

$$\rho = M/A = \frac{M}{\frac{1}{2}(a^2 - b^2)} \Rightarrow F = \frac{-4 G m M \ln \left( \frac{a}{b} \right)}{(a^2 - b^2)}$$

4. Find the force of attraction of a thin uniform rod of length  $2a$  on a particle of mass  $m$  placed at a distance  $b$  from its midpoint.

**Solution** : Let's begin with the following graphic :



The gravitational force on a test mass at  $m$  due to an element of the line denoted by  $dx$  is:

$$dF = -\frac{G m dM}{r^2} = -\frac{G m dM}{x^2 + b^2}$$

where  $m$  is the mass of the test particle and  $dM$  is the mass of the element  $dx$ . If the bar is uniform, its mass can be written as  $dM = \rho dx$  and we have:

$$dF = \frac{-G m \rho dx}{x^2 + b^2}$$

We can use symmetry arguments again to show the  $x$  components of force will cancel, and the  $y$  component of force is given by:

$$dF_y = - \frac{-G m \rho \cos \theta dx}{x^2 + b^2}$$

and the geometry of the situation shows us that

$$\cos \theta = \frac{b}{\sqrt{x^2 + b^2}}$$

Thus, the total force at  $b$  can be found by integrating from  $-a$  to  $a$  :

$$F_y = -G m \rho b \int_{-a}^a \frac{dx}{(x^2 + b^2)^{3/2}}$$

and symmetry also allows us to write :

$$F_y = -2 G m \rho b \int_0^a \frac{dx}{(x^2 + b^2)^{3/2}}$$

We can solve this by making the substitution  $x = b \tan \theta$ , then

$$x^2 + b^2 = b^2 \tan^2 \theta + b^2 = b^2 \sec^2 \theta$$

and

$$dx = b \sec^2 \theta d\theta$$

so that our integral becomes:

$$\begin{aligned} F_y &= -2 G m \rho b \int_0^{\tan^{-1}(a/b)} \frac{b \sec^2 \theta d\theta}{(b^2 \sec^2 \theta)^{3/2}} = \frac{-2 G m \rho}{b} \int_0^{\tan^{-1}(a/b)} \cos \theta d\theta \\ &= -\frac{2 G m \rho}{b} \sin \theta = \frac{-2 G m \rho a}{b \sqrt{a^2 + b^2}} \end{aligned}$$

and we can eliminate  $\rho$  by setting  $\rho = M/2a$  so that:

$$F_y = \frac{-G m M}{b \sqrt{a^2 + b^2}}$$